Communications in Contemporary Mathematics
Vol. 21, No. 1 (2019) 1850017 (13 pages)

DOI: 10.1142/S0219199718500177

On anisotropic Sobolev spaces

Hoai-Minh Nguyen
Department of Mathematics, EPFL SB CAMA
Station 8 CH-1015 Lausanne, Switzerland
hoai-minh.nguyen@epfl.ch

Marco Squassina
Dipartimento di Matematica e Fisica
Università Cattolica del Sacro Cuore
Via dei Musei 41, I-25121 Brescia, Italy
marco.squassina@unicatt.it

Received 2 October 2017
Accepted 19 February 2018
Published 15 May 2018

We investigate two types of characterizations for anisotropic Sobolev and BV spaces. In particular, we establish anisotropic versions of the Bourgain–Brezis–Mironescu formula, including the magnetic case both for Sobolev and BV functions.

Keywords: Nonlocal functionals; characterization of anisotropic Sobolev spaces.

Mathematics Subject Classification 2010: 46E35, 28D20, 82B10, 49A50

1. Introduction and Results

1.1. *Overview*

Around 2001, Bourgain, Brezis and Mironescu, investigated (cf. [2, 3, 6]) the asymptotic behavior of a class on nonlocal functionals on a domain $\Omega \subset \mathbb{R}^N$, including those related to the norms of the fractional Sobolev spaces $W^{s,p}(\mathbb{R}^N)$ as $s \nearrow 1$. In the case $\Omega = \mathbb{R}^N$, their later result can be formulated as follows: if $p > 1$ and $u \in W^{1,p}(\mathbb{R}^N)$, then

$$\lim_{s \nearrow 1}(1-s)\iint_{\mathbb{R}^{2N}} \frac{|u(x)-u(y)|^p}{|x-y|^{N+ps}}dx\,dy = K_{p,N}\int_{\mathbb{R}^N}|\nabla u|^p dx, \tag{1.1}$$

where

$$K_{p,N} = \frac{1}{p}\int_{\mathbb{S}^{N-1}}|\boldsymbol{\omega}\cdot x|^p d\sigma, \tag{1.2}$$

being $\boldsymbol{\omega} \in \mathbb{S}^{N-1}$ any fixed vector. Here and in what follows, for a vector $x \in \mathbb{R}^N$, $|x|$ denotes its Euclidean norm.

Given a convex, symmetric subset $K \subset \mathbb{R}^N$ containing the origin, let $\|\cdot\|_K$ be the norm in $\mathbb{R}^N$ which admits as unit ball the set K, i.e.

$$\|x\|_K := \inf\Big\{\lambda > 0 : \frac{x}{\lambda} \in K\Big\}. \tag{1.3}$$

It is rather natural to wonder what happens to formula (1.1) by replacing in the singular kernel $|x-y|$ with its anisotropic version $\|x-y\|_K$. In 2014, Ludwig [18, 19] proved that, for a compactly supported function $u \in W^{1,p}(\mathbb{R}^N)$, there holds

$$\lim_{s\nearrow 1}(1-s)\iint_{\mathbb{R}^{2N}} \frac{|u(x)-u(y)|^p}{\|x-y\|_K^{N+ps}}dx\,dy = \int_{\mathbb{R}^N} \|\nabla u\|^p_{Z_p^*K}dx. \tag{1.4}$$

Here $\|\cdot\|_{Z_p^*K}$ is the norm associated with the convex set Z_p^*K which is the polar L_p moment body of K (see (1.6) and (1.7)); such quantities were involved in recent important applications within convex geometry and probability theory, see e.g. [14, 15, 17] and the references therein. Thus, changing the norm in the nonlocal functional produces anisotropic effects in the singular limit. The norm

$$v \mapsto \|v\|_{Z_p^*K},$$

can be explicitly written and, in the particular case $\|\cdot\|_K = |\cdot|$ (Euclidean case), then $K = B_1$, the unit ball of $\mathbb{R}^N$, and the results are consistent with classical formulas, since $\|\cdot\|_{Z_p^*B} = \sqrt[p]{K_{p,N}}\,|\cdot|$.

Ludwig's proof of formula (1.4) relies on a reduction argument involving the one-dimensional version of the Bourgain–Brezis–Mironescu formula in the Euclidean setting jointly with the *Blaschke-Petkantschin* geometric integration formula (cf. [27, Theorem 7.2.7]), namely

$$\int_{\mathbb{R}^N}\int_{\mathbb{R}^N} f(x,y)dxdy = \int_{\mathrm{Aff(N,1)}}\int_L\int_L f|_{L\times L}(x,y)|x-y|^{N-1}d\mathscr{H}^1(x)d\mathscr{H}^1(y)dL,$$

where $\mathscr{H}^1$ is the one-dimensional Hausdorff measure on $\mathbb{R}^N$, $\mathrm{Aff}(N,1)$ is the affine Grassmannian of lines in $\mathbb{R}^N$ and dL denotes the integration with respect to a Haar measure on $\mathrm{Aff}(N,1)$.

Around 2006, motivated by an estimate for the topological degree raising in the framework of Ginzburg–Landau equations [4], a new alternative characterization of the Sobolev spaces was introduced (cf. [5, 20, 21]). As a result, for every $u \in W^{1,p}(\mathbb{R}^N)$ with $p > 1$, there holds

$$\lim_{\delta\to 0}\iint_{\{|u(y)-u(x)|>\delta\}} \frac{\delta^p}{|x-y|^{N+p}}dx\,dy = K_{p,N}\int_{\mathbb{R}^N}|\nabla u|^p dx, \tag{1.5}$$

where $K_{p,N}$ is the constant appearing in (1.2). It is thus natural to wonder if, replacing $|x-y|$ in the singular kernel with the corresponding anisotropic version $\|x-y\|_K$, produces in the limit the same result as in formula (1.4).

The previous two characterizations were also considered for $p = 1$. BV functions are involved in this case, see [2, 5, 12, 20]. Other properties related to these characterizations can be found in [8, 9, 11, 22–24, 26]. Both the characterizations (for

the Euclidean norm) were recently extended to the case of *magnetic Sobolev* and *BV spaces* [24, 25, 28]. More general nonlocal functionals have been investigated in [7–11].

1.2. *Anisotropic spaces*

In this section, we introduce anisotropic magnetic Sobolev and BV spaces. For this end, complex numbers and notations are involved. Let $p \geq 1$ and consider the complex space $(\mathbb{C}^N, |\cdot|_p)$ endowed with

$$|z|_p := (|(\Re z_1, \ldots, \Re z_N)|^p + |(\Im z_1, \ldots, \Im z_N)|^p)^{1/p},$$

where $\Re a$ and $\Im a$ denote the real and imaginary parts of $a \in \mathbb{C}$. Recall that $|x|$ is the *Euclidean* norm of $x \in \mathbb{R}^N$. Note that $|z|_p = |z|$ for $z \in \mathbb{R}^N$. Let $\|\cdot\|_K$ be the norm as in (1.3). We set

$$\|v\|_{Z_p^* K} := \left(\frac{N+p}{p}\int_K |v \cdot x|_p^p dx\right)^{1/p}, \quad \text{for } v \in \mathbb{C}^N. \tag{1.6}$$

The set $Z_p^* K \subset \mathbb{C}^N$ which is defined as

$$Z_p^* K := \{v \in \mathbb{C}^N : \|v\|_{Z_p^* K} \leq 1\} \tag{1.7}$$

is called the (complex) polar L_p-moment body of K. Denote $L^p(\mathbb{R}^N, \mathbb{C})$ the Lebesgue space of functions $u : \mathbb{R}^N \to \mathbb{C}$ such that

$$\|u\|_{L^p(\mathbb{R}^N)} := \left(\int_{\mathbb{R}^N} |u|_p^p dx\right)^{1/p} < \infty.$$

For a locally bounded function $A : \mathbb{R}^N \to \mathbb{R}^N$ (magnetic potential), set

$$[u]_{W^{1,p}_{A,K}(\mathbb{R}^N)} := \left(\int_{\mathbb{R}^N} \|\nabla u - \mathrm{i}A(x)u\|_{Z_p^* K}^p dx\right)^{1/p}.$$

Let $W^{1,p}_{A,K}(\mathbb{R}^N)$ be the space of $u \in L^p(\mathbb{R}^N, \mathbb{C})$ such that $[u]_{W^{1,p}_{A,K}(\mathbb{R}^N)} < \infty$ with the norm

$$\|u\|_{W^{1,p}_{A,K}(\mathbb{R}^N)} := \left(\|u\|^p_{L^p(\mathbb{R}^N)} + [u]^p_{W^{1,p}_{A,K}(\mathbb{R}^N)}\right)^{1/p}.$$

Denote $\|\cdot\|_{Z_1^* K^*}$ the dual norm of the norm $\|\cdot\|_{Z_1^* K}$ on $\mathbb{R}^N$, namely for $v \in \mathbb{R}^N$

$$\|v\|_{Z_1^* K^*} := \sup\{\langle v, w\rangle_{\mathbb{R}^N} : w \in \mathbb{R}^N,\ \|w\|_{Z_1^* K} \leq 1\}, \quad \text{with}$$

$$\langle v, w\rangle_{\mathbb{R}^N} = \sum_{j=1}^N v_j w_j, \quad \forall\, v, w \in \mathbb{R}^N.$$

For a complex function $u \in L^1_{\mathrm{loc}}(\mathbb{R}^N)$, as in [25], we define

$$|Du|_{A,K} := C_{1,A,K,u} + C_{2,A,K,u},$$

where

$$C_{1,A,K,u} := \sup\Big\{\int_{\mathbb{R}^N} \Re u \operatorname{div}\varphi - A\cdot\varphi \Im u\,dx,\ \varphi\in C_c^1(\mathbb{R}^N,\mathbb{R}^N) \text{ with } \|\varphi(x)\|_{Z_1^*K^*}\le 1 \text{ in } \mathbb{R}^N\Big\},$$

$$C_{2,A,K,u} := \sup\Big\{\int_{\mathbb{R}^N} \Im u \operatorname{div}\varphi + A\cdot\varphi \mathbb{R} u\,dx,\ \varphi\in C_c^1(\mathbb{R}^N,\mathbb{R}^N) \text{ with } \|\varphi(x)\|_{Z_1^*K^*}\le 1 \text{ in } \mathbb{R}^N\Big\}.$$

We say that $u\in BV_{A,K}(\mathbb{R}^N)$ if $u\in L^1(\mathbb{R}^N)$ and $|Du|_{A,K}<\infty$ and in this case we formally set

$$|Du|_{A,K} = \int_{\mathbb{R}^N} \|\nabla u - \mathrm{i}A(x)u\|_{Z_1^*K}\,dx. \tag{1.8}$$

The space $BV_{A,K}(\mathbb{R}^N)$ is a Banach space [25] equipped the norm

$$\|u\|_{A,K} = \|u\|_{L^1(\mathbb{R}^N)} + |Du|_{A,K}, \quad u\in BV_{A,K}(\mathbb{R}^N).$$

1.3. *Main results*

The goal of this paper is to extend the two characterizations mentioned above to anisotropic magnetic Sobolev and BV spaces. Our approach is in the spirit of the works on the Euclidean spaces. In particular, we make no use of the Blaschke-Petkantschin geometric integration formula as in the work of Ludwig.

Let $A:\mathbb{R}^N\to\mathbb{R}^N$ be measurable and locally bounded. Set

$$\Psi_u(x,y) := e^{\mathrm{i}(x-y)\cdot A\left(\frac{x+y}{2}\right)}u(y), \quad x,y\in\mathbb{R}^N.$$

Motivated by the study of the interaction of particles in the presence of a magnetic field, see e.g., [1, 16] and references therein, Ichinose [16] considered the nonlocal functional

$$H_A^s(\mathbb{R}^N)\ni u\mapsto \iint_{\mathbb{R}^{2N}} \frac{|u(x)-e^{\mathrm{i}(x-y)\cdot A(\frac{x+y}{2})}u(y)|^2}{|x-y|^{N+2s}}dx\,dy,$$

for $s\in(0,1)$, and established that its gradient is the fractional Laplacian associated with the magnetic field A via a probabilistic argument. As in the spirit of the previous results, the quantity Ψ_u has been recently involved in the characterization of magnetic Sobolev and BV functions. In this paper, we establish the following anisotropic magnetic version of (1.5).

Theorem 1.1. *Let $p>1$ and let $A:\mathbb{R}^N\to\mathbb{R}^N$ be Lipschitz. Then, for every $u\in W^{1,p}_{A,K}(\mathbb{R}^N)$,*

$$\lim_{\delta\searrow 0}\int_{\mathbb{R}^N}\int_{\mathbb{R}^N} \frac{\delta^p}{\|x-y\|_K^{N+p}}\mathbf{1}_{\{|\Psi_u(x,y)-\Psi_u(x,x)|_p>\delta\}}dx\,dy = \int_{\mathbb{R}^N}\|\nabla u-\mathrm{i}A(x)u\|^p_{Z_p^*K}dx.$$

If $p = 1$, one can show (see Remark 2.1) that, for $u \in W^{1,1}_{A,K}(\mathbb{R}^N)$,

$$\lim_{\delta \searrow 0} \int_{\mathbb{R}^N} \int_{\mathbb{R}^N} \frac{\delta}{\|x-y\|_K^{N+1}} \mathbf{1}_{\{|\Psi_u(x,y)-\Psi_u(x,x)|_p > \delta\}} dx\, dy$$
$$\geq \int_{\mathbb{R}^N} \|\nabla u - \mathrm{i}A(x)u\|_{Z_1^* K} dx.$$

Nevertheless, such an inequality does not hold in general for $u \in BV_A(\mathbb{R}^N)$ even in the case where $A \equiv 0$ and K is the unit ball (see [11, Pathology 3]). In the case $A = 0$, one has

$$\lim_{\delta \searrow 0} \int_{\mathbb{R}^N} \int_{\mathbb{R}^N} \frac{\delta}{\|x-y\|_K^{N+1}} \mathbf{1}_{|u(y)-u(x)|>\delta\}} dx\, dy \geq C \int_{\mathbb{R}^N} \|\nabla u\|_{Z_1^* K} dx,$$

for some positive constant $0 < C < K_{N,1}$. This inequality is a direct consequence of the corresponding result in the Euclidean setting in [5].

We next discuss the BBM formula for the anisotropic magnetic setting. Let (ρ_n) be a sequence of non-negative radial mollifiers such that

$$\lim_{n \to +\infty} \int_\delta^\infty \rho_n(r) r^{N-1-p} dr = 0, \quad \text{for all } \delta > 0 \quad \text{and}$$
$$\int_0^1 \rho_n(r) r^{N-1} dr = 1. \tag{1.9}$$

Here is the anisotropic magnetic BBM formula.

Theorem 1.2. *Let $p \geq 1$, let $A : \mathbb{R}^N \to \mathbb{R}^N$ be Lipschitz, and let $\{\rho_n\}_{n\in\mathbb{N}}$ be a sequence of non-negative radial mollifiers satisfying* (1.9)*. Then, for $u \in W^{1,p}_{A,K}(\mathbb{R}^N)$,*

$$\lim_{n \to +\infty} \iint_{\mathbb{R}^{2N}} \frac{|\Psi_u(x,y) - \Psi_u(x,x)|^p}{\|x-y\|_K^p} \rho_n(\|x-y\|_K) dx\, dy$$
$$= p \int_{\mathbb{R}^N} \|\nabla u - \mathrm{i}A(x)u\|^p_{Z_p^* K} dx.$$

Furthermore, if $p = 1$ and $u \in BV_{A,K}(\mathbb{R}^N)$ the formula holds with the agreement (1.8)*.*

Remark 1.3. Let (s_n) be a positive sequence converging to 0 and set, for $n \geq 1$,

$$\rho_n(r) = \frac{p(1-s_n)}{r^{N+ps_n-p}}, \quad r > 0.$$

Then (ρ_n) satisfy (1.9). Applying Theorem 1.2, one rediscovers the results of Ludwig.

Remark 1.4. Theorems 1.1 and 1.2 provide the full solution of a problem arised by Giuseppe Mingione on September 21th, 2016, at the end of the seminar "*Another triumph for De Giorgi's Gamma convergence*" by Haim Brezis at the conference "*A Mathematical tribute to Ennio De Giorgi*", held in Pisa from 19th to 23th September 2016.

The above results provide an extension of [2, 5, 12, 18–20, 24–26, 28] to the anisotropic case.

2. Proof of Theorem 1.1

Set, for $p \geq 1$,

$$I_\delta^K(u) := \int_{\mathbb{R}^N}\int_{\mathbb{R}^N} \frac{\delta^p}{\|x-y\|_K^{N+p}} \mathbf{1}_{\{|\Psi_u(x,y)-\Psi_u(x,x)|_p>\delta\}} dx\, dy, \quad \text{for } u \in L^1_{\mathrm{loc}}(\mathbb{R}^N).$$

It is clear that, for $u, v \in W^{1,p}_{A,K}(\mathbb{R}^N)$ and $0 < \varepsilon < 1$,

$$I_\delta^K(u) \leq (1-\varepsilon)^{-p} I^K_{(1-\varepsilon)\delta}(v) + \varepsilon^{-p} I^K_{\varepsilon\delta}(u-v). \tag{2.1}$$

Applying [24, Theorem 3.1], we have, for $p > 1$ and $u \in W^{1,p}_{A,K}(\mathbb{R}^N)$,

$$I_\delta^K(u) \leq C_{N,p,K}\left(\int_{\mathbb{R}^N} |\nabla u - \mathrm{i}A(x)u|_p^p dx + (\|\nabla A\|^p_{L^\infty(\mathbb{R}^N)} + 1)\int_{\mathbb{R}^N} |u|_p^p dx\right),$$

for some positive constant $C_{N,p,K}$ depending only on N, p, and K. By the density of $C^1_c(\mathbb{R}^N)$ in $W^{1,p}_A(\mathbb{R}^N)$, it hence suffices to consider the case $u \in C^1_c(\mathbb{R}^N)$ which will be assumed from later on.

By a change of variables, we have

$$\begin{aligned}&\int_{\mathbb{R}^N}\int_{\mathbb{R}^N} \frac{\delta^p}{\|x-y\|_K^{N+p}} \mathbf{1}_{\{|\Psi_u(x,y)-\Psi_u(x,x)|_p>\delta\}} dx\, dy \\ &\quad = \int_{\mathbb{R}^N}\int_{\mathbb{S}^{N-1}}\int_0^\infty \frac{1}{\|\sigma\|_K^{N+p} h^{1+p}} \mathbf{1}_{\{|\Psi_u(x,x+\delta h\sigma)-\Psi_u(x,x)|_p>\delta\}}\, dh\, d\sigma\, dx.\end{aligned}$$

Using the fact

$$\lim_{\delta\to 0} \frac{|\Psi_u(x, x+\delta h\sigma) - \Psi_u(x,x)|_p}{\delta} = |(\nabla u - \mathrm{i}A(x)u)\cdot\sigma|_p h, \tag{2.2}$$

as in the proof of [24, Lemma 3.3], we obtain

$$\begin{aligned}&\lim_{\delta\to 0}\int_{\mathbb{R}^N}\int_{\mathbb{R}^N} \frac{\delta^p}{\|x-y\|_K^{N+p}} \mathbf{1}_{\{|\Psi_u(x,y)-\Psi_u(x,x)|_p>\delta\}} \\ &\quad = \frac{1}{p}\int_{\mathbb{R}^N}\int_{\mathbb{S}^{N-1}} \frac{|(\nabla u - \mathrm{i}A(x)u)\cdot\sigma|_p^p}{\|\sigma\|_K^{N+p}} d\sigma\, dx.\end{aligned} \tag{2.3}$$

Since we have

$$\begin{aligned}(N+p)\int_K |v\cdot y|_p^p dy &= (N+p)\int_{\mathbb{S}^{N-1}}\int_0^{1/\|\sigma\|_K} |v\cdot\sigma|_p^p t^{N-1+p} dt\, d\sigma \\ &= \int_{\mathbb{S}^{N-1}} \frac{|v\cdot\sigma|_p^p}{\|\sigma\|_K^{N+p}} d\sigma,\end{aligned} \tag{2.4}$$

the assertion follows.

Remark 2.1. In the case $u \in W^{1,1}_A(\mathbb{R}^N)$, by Fatou's lemma, as in (2.3), one has, with $p = 1$,

$$\lim_{\delta\to 0} I^K_\delta(u) \geq \int_{\mathbb{R}^N}\int_{\mathbb{S}^{N-1}}\int_0^\infty_{\{|(\nabla u(x)-\mathrm{i}A(x)u(x))\cdot\sigma|_1 h>1\}} \frac{1}{\|\sigma\|_K^{N+p}h^2} dh\, d\sigma\, dx$$
$$= \frac{1}{p}\int_{\mathbb{R}^N}\int_{\mathbb{S}^{N-1}} \frac{|(\nabla u - \mathrm{i}A(x)u)\cdot\sigma|_1}{\|\sigma\|_K^{N+1}} d\sigma\, dx.$$

This implies

$$\lim_{\delta\searrow 0}\int_{\mathbb{R}^N}\int_{\mathbb{R}^N} \frac{\delta}{\|x-y\|_K^{N+1}} \mathbf{1}_{\{|\Psi_u(x,y)-\Psi_u(x,x)|_1>\delta\}} dx\, dy \geq \int_{\mathbb{R}^N} \|\nabla u - \mathrm{i}A(x)u\|_{Z_1^* K} dx.$$

3. Proof of Theorem 1.2

3.1. *Proof of Theorem* 1.2 *for* $p > 1$

Using [24, Theorem 2.1] without loss of generality, one might assume that $u \in C^1_c(\mathbb{R}^N)$. Note that

$$\iint_{\mathbb{R}^{2N}} \frac{|\Psi_u(x,y)-\Psi_u(x,x)|_p^p}{\|x-y\|_K^p}\rho_n(\|x-y\|_K)dx\, dy$$
$$= \int_{\mathbb{R}^N}\int_{\mathbb{S}^{N-1}}\int_0^\infty \frac{|\Psi_u(x,x+h\sigma)-\Psi_u(x,x)|_p^p}{\|\sigma\|_K^p h^p}\rho_n(\|\sigma\|_K h)h^{N-1}dh\, d\sigma\, dx.$$

Using (2.2), one then can check that, for $p \geq 1$ and $u \in C^1_c(\mathbb{R}^N)$,

$$\lim_{n\to+\infty}\iint_{\{|x-y|\leq 1\}} \frac{|\Psi_u(x,y)-\Psi_u(x,x)|_p^p}{\|x-y\|_K^p}\rho_n(\|x-y\|_K)dx\, dy$$
$$= \int_{\mathbb{R}^N}\int_{\mathbb{S}^{N-1}} \frac{|(\nabla u - \mathrm{i}A(x)u)\cdot\sigma|_p^p}{\|\sigma\|_K^p} d\sigma\, dx \lim_{n\to+\infty}\int_0^1 \rho_n(\|\sigma\|_K h)h^{N-1}dh.$$

Furthermore, observe that

$$\iint_{\{|x-y|>1\}} \frac{|\Psi_u(x,y)-\Psi_u(x,x)|_p^p}{\|x-y\|_K^p}\rho_n(\|x-y\|_K)dx\, dy$$
$$\leq C\|u\|_{L^p}^p \int_1^\infty h^{N-1-p}\rho_n(\|\sigma\|_K h)dh.$$

Therefore, for $p \geq 1$ and $u \in C^1_c(\mathbb{R}^N)$, on account of (1.9), we obtain

$$\lim_{n\to+\infty}\iint_{\mathbb{R}^{2N}} \frac{|\Psi_u(x,y)-\Psi_u(x,x)|_p^p}{\|x-y\|_K^p}\rho_n(\|x-y\|_K)dx\, dy$$
$$= \int_{\mathbb{R}^N}\int_{\mathbb{S}^{N-1}} \frac{|(\nabla u - \mathrm{i}A(x)u)\cdot\sigma|_p^p}{\|\sigma\|_K^{N+p}} d\sigma\, dx. \tag{3.1}$$

The conclusion now follows from (2.4).

3.2. *Proof of Theorem* 1.2 *for* $p = 1$

We first present some preliminary results. The first one is the following lemma.

Lemma 3.1. *Let* $u \in W^{1,1}_{A,K}(\mathbb{R}^N)$. *Then*

$$|Du|_{A,K} = \int_{\mathbb{R}^N} \|\nabla u - \mathrm{i}A(x)u\|_{Z_1^* K} dx.$$

Proof. The proof is quite standard and based on integration by parts after noting that

$$\|\nabla u - \mathrm{i}A(x)u\|_{Z_1^* K} = \|\nabla \Re u - A(x)\Im u\|_{Z_1^* K} + \|\nabla \Im u + A(x)\Re u\|_{Z_1^* K},$$

since $A(x) \in \mathbb{R}^N$ for $x \in \mathbb{R}^N$. The details are left to the reader. □

Lemma 3.2. *Let* $u \in BV_A(\mathbb{R}^N)$ *and* $(u_n) \subset BV_A(\mathbb{R}^N)$. *Assume that*

$$\lim_{n\to+\infty} u_n = u \quad \text{in } L^1(\mathbb{R}^N).$$

Then

$$\liminf_{n\to+\infty} |Du_n|_{A,K} \geq |Du|_{A,K}.$$

Proof. One can check that

$$\liminf_{n\to+\infty} C_{1,A,K,u_n} \geq C_{1,A,K,u} \quad \text{and} \quad \liminf_{n\to+\infty} C_{2,A,K,u_n} \geq C_{2,A,K,u}.$$

The conclusion follows. □

For $r > 0$, let B_r denote the ball centered at the origin and of radius r. We have

Lemma 3.3. *Let* $u \in BV_A(\mathbb{R}^N)$ *and let* (τ_m) *be a sequence of non-negative mollifiers with* $\operatorname{supp} \tau_m \subset B_{1/m}$ *which is normalized by the condition* $\int_{\mathbb{R}^N} \tau_m(x)dx = 1$. *Set* $u_m = \tau_m * u$. *Assume that* A *is Lipschitz. Then*

$$\lim_{m\to+\infty} |Du_m|_{A,K} = |Du|_{A,K}.$$

Proof. The proof is quite standard, see e.g., [13] and also [25]. Let $\varphi \in C^1_c(\mathbb{R}^N)$ be such that

$$\|\varphi(x)\|_{Z_1^* K^*} \leq 1 \quad \text{in } \mathbb{R}^N.$$

We have

$$\begin{aligned}
&\int_{\mathbb{R}^N} \Re u_m \operatorname{div}\varphi - A \cdot \varphi \Im u_m \, dx \\
&\quad = \int_{\mathbb{R}^N} \Re u \operatorname{div}\varphi_m - A \cdot \varphi_m \Im u \, dx \\
&\qquad + \int_{\mathbb{R}^N} \int_{\mathbb{R}^N} (A(x) - A(x-y)) \cdot \varphi(x-y)\tau_m(y)u(x)dx\, dy. \qquad (3.2)
\end{aligned}$$

Since

$$\|\varphi_m(x)\|_{Z_1^* K^*} \leq \sup_y \|\varphi(y)\|_{Z_1^* K^*} \leq 1,$$

we have

$$\left| \int_{\mathbb{R}^N} \Re u \operatorname{div} \varphi_m - A \cdot \varphi_m \Im u \, dx \right| \leq C_{1,A,K,u}. \tag{3.3}$$

Since $\operatorname{supp} \tau_m \subset B_{1/m}$, one can check that

$$\left| \int_{\mathbb{R}^N} \int_{\mathbb{R}^N} (A(x) - A(x-y)) \cdot \varphi(x-y) \tau_m(y) u(x) dx \, dy \right| \leq \frac{C \|\nabla A\|_{L^\infty} \|u\|_{L^1}}{m}. \tag{3.4}$$

A combination of (3.2), (3.3), and (3.4) yields

$$\limsup_{m \to +\infty} C_{1,A,K,u_m} \leq C_{1,A,K,u}.$$

Similarly, we obtain

$$\limsup_{m \to +\infty} C_{2,A,K,u_m} \leq C_{2,A,K,u}$$

and the conclusion follows from Lemma 3.2. □

We are ready to give the following proof of theorem.

Proof of Theorem 1.2 for $p = 1$. Let (τ_m) be a sequence of non-negative mollifiers with $\operatorname{supp} \tau_m \subset B_{1/m}$ which is normalized by the condition $\int_{\mathbb{R}^N} \tau_m(x) \, dx = 1$. Set $u_m = u * \tau_m$. As in the proof of [24, Lemma 2.4], we have

$$\begin{aligned} &\iint_{\mathbb{R}^{2N}} \frac{|\Psi_{u_m}(x,y) - \Psi_{u_m}(x,x)|_1}{\|x-y\|_K} \rho_n(\|x-y\|_K) dx \, dy \\ &\quad \leq \iint_{\mathbb{R}^{2N}} \frac{|\Psi_u(x,y) - \Psi_u(x,x)|_1}{\|x-y\|_K} \rho_n(\|x-y\|_K) dx \, dy \\ &\qquad + C \int_{\mathbb{R}^N} \int_{\mathbb{R}^N} \int_{\mathbb{R}^N} |z| \tau_m(z) \rho_n(\|x-y\|_K) u(y) dz \, dx \, dy. \end{aligned}$$

We have

$$\begin{aligned} &\lim_{m \to +\infty} \lim_{n \to +\infty} \iint_{\mathbb{R}^{2N}} \frac{|\Psi_{u_m}(x,y) - \Psi_{u_m}(x,x)|_1}{\|x-y\|_K} \rho_n(\|x-y\|_K) dx \, dy \\ &\quad \geq \lim_{m \to +\infty} \int_{\mathbb{R}^N} \|\nabla u_m - \mathrm{i} A(x) u_m\|_{Z_1^* K} \, dx \overset{\text{Lemma 3.3}}{=} \int_{\mathbb{R}^N} \|\nabla u - \mathrm{i} A(x) u\|_{Z_1^* K} \, dx \end{aligned}$$

and, since $\operatorname{supp} \tau_m \subset B_{1/m}$,

$$\int_{\mathbb{R}^N} \int_{\mathbb{R}^N} \int_{\mathbb{R}^N} |z| \tau_m(z) \rho_n(\|x-y\|_K) u(y) dz \, dx \, dy \leq \frac{C}{m}.$$

It follows that

$$\liminf_{n\to+\infty}\iint_{\mathbb{R}^{2N}}\frac{|\Psi_u(x,y)-\Psi_u(x,x)|_1}{\|x-y\|_K}\rho_n(\|x-y\|_K)dx\,dy$$
$$\geq\int_{\mathbb{R}^N}\|\nabla u-\mathrm{i}A(x)u\|_{Z_1^*K}dx.$$

We also have, by Fatou's lemma,

$$\iint_{\mathbb{R}^{2N}}\frac{|\Psi_u(x,y)-\Psi_u(x,x)|_1}{\|x-y\|_K}\rho_n(\|x-y\|_K)dx\,dy$$
$$\leq\liminf_{m\to+\infty}\iint_{\mathbb{R}^{2N}}\frac{|\Psi_{u_m}(x,y)-\Psi_{u_m}(x,x)|_1}{\|x-y\|_K}\rho_n(\|x-y\|_K)dx\,dy. \tag{3.5}$$

We next derive an upper bound for the RHS of (3.5). Let $v\in W^{1,1}_{A,K}(\mathbb{R}^N)\cap C^\infty(\mathbb{R}^N)$. We have

$$\iint_{\mathbb{R}^{2N}}\frac{|\Psi_v(x,y)-\Psi_v(x,x)|_1}{\|x-y\|_K}\rho_n(\|x-y\|_K)dx\,dy$$
$$=\int_{\mathbb{R}^N}\int_{\mathbb{S}^{N-1}}\int_0^\infty\frac{|\Psi_v(x,x+h\sigma)-\Psi_v(x,x)|_1}{\|h\sigma\|_K}\rho_n$$
$$\times(h\|\sigma\|_K)h^{N-1}dh\,d\sigma\,dx. \tag{3.6}$$

Using the fact

$$\frac{\partial\Psi_v(x,y)}{\partial y}=e^{\mathrm{i}(x-y)\cdot A(\frac{x+y}{2})}\nabla v(y)-\mathrm{i}\left\{A\left(\frac{x+y}{2}\right)+\frac{1}{2}(y-x)\cdot\nabla A\left(\frac{x+y}{2}\right)\right\}$$
$$\times e^{\mathrm{i}(x-y)\cdot A(\frac{x+y}{2})}v(y),$$

and applying the mean value theorem, we obtain

$$\iint_{\mathbb{R}^{2N}}\frac{|\Psi_v(x,y)-\Psi_v(x,x)|_1}{\|x-y\|_K}\rho_n(\|x-y\|_K)dx\,dy$$
$$\leq\int_{\mathbb{R}^N}\int_{\mathbb{S}^{N-1}}\int_0^1\int_0^1|\big(\nabla v-\mathrm{i}Av\big)\cdot\sigma|_1(x+th\sigma)$$
$$\times\frac{1}{\|\sigma\|_K}h^{N-1}\rho_n(h\|\sigma\|_K)\,dt\,dh\,d\sigma\,dx$$
$$+C\int_{\mathbb{R}^N}\int_{\mathbb{S}^{N-1}}\int_0^1\int_0^1|v(x+th\sigma)|\frac{\|\nabla A\|_{L^\infty}h}{\|\sigma\|_K}h^{N-1}\rho_n(h\|\sigma\|_K)\,dt\,dh\,d\sigma\,dx$$
$$+C\int_{\substack{\mathbb{R}^N\\ \{|x-y|>1\}}}\int_{\mathbb{R}^N}(|v(x)|+|v(y)|)\frac{1}{\|x-y\|_K}\rho_n(\|x-y\|_K)\,dx\,dy. \tag{3.7}$$

One can check that

$$\int_{\mathbb{R}^N}\int_{\mathbb{S}^{N-1}}\int_0^1\int_0^1 |\big(\nabla v - \mathrm{i}Av\big)\cdot\sigma|_1\big(x+th\sigma\big)\frac{1}{\|\sigma\|_K}h^{N-1}\rho_n(h\|\sigma\|_K)dt\,dh\,d\sigma\,dx$$
$$\leq \int_{\mathbb{R}^N}\int_{\mathbb{S}^{N-1}}\frac{|(\nabla v - \mathrm{i}A(x)v)\cdot\sigma|_1}{\|\sigma\|_K^{N+1}}\,d\sigma\,dx\int_0^\lambda h^{N-1}\rho_n(h)dh, \tag{3.8}$$

and

$$\int_{\mathbb{R}^N}\int_{\mathbb{S}^{N-1}}\int_0^1\int_0^1 |v(x+th\sigma)|\frac{\|\nabla A\|_{L^\infty}h}{\|\sigma\|_K}h^{N-1}\rho_n(h\|\sigma\|_K)dt\,dh\,d\sigma\,dx$$
$$\leq C_K\|\nabla A\|_{L^\infty}\|v\|_{L^1}\int_0^\lambda h^N\rho_n(h)dh, \tag{3.9}$$

where $\lambda = \max\{\|\sigma\|_K : \sigma\in\mathbb{S}^{N-1}\}$. A combination of (3.6)–(3.9) yields

$$\iint_{\mathbb{R}^{2N}}\frac{|\Psi_v(x,y)-\Psi_v(x,x)|_1}{\|x-y\|_K}\rho_n(\|x-y\|_K)dx\,dy$$
$$\leq \int_0^\lambda h^{N-1}\rho_n(h)dh\int_{\mathbb{R}^N}\|\nabla v - \mathrm{i}A(x)v\|_{Z_1^*K}dx$$
$$+C_K(\|\nabla A\|_{L^\infty}+1)\|v\|_{L^1}\left(\int_0^\lambda h^N\rho_n(h)\,dh + \int_1^\infty h^{N-2}\rho_n(h)dh\right). \tag{3.10}$$

Using Lemma 3.3, we derive from (3.5) and (3.10) that

$$\iint_{\mathbb{R}^{2N}}\frac{|\Psi_u(x,y)-\Psi_u(x,x)|_1}{\|x-y\|_K}\rho_n(\|x-y\|_K)dx\,dy$$
$$\leq \int_0^\lambda h^{N-1}\rho_n(h)\int_{\mathbb{R}^N}\|\nabla u - \mathrm{i}A(x)u\|_{Z_1^*K}dx$$
$$+C_K\|\nabla A\|_{L^\infty}\|u\|_{L^1}\left(\int_0^\lambda h^N\rho_n(h)dh + \int_1^\infty h^{N-2}\rho_n(h)dh\right),$$

which yields, by (1.9),

$$\limsup_{n\to+\infty}\iint_{\mathbb{R}^{2N}}\frac{|\Psi_u(x,y)-\Psi_u(x,x)|_1}{\|x-y\|_K}\rho_n(\|x-y\|_K)dx\,dy$$
$$\leq \int_{\mathbb{R}^N}\|\nabla v - \mathrm{i}A(x)v\|_{Z_1^*K}dx.$$

The proof is complete. □

Acknowledgment

The second author is member of *Gruppo Nazionale per l'Analisi Matematica, la Probabilità e le loro Applicazioni* (GNAMPA) of the *Istituto Nazionale di Alta*

Matematica (INdAM). Part of this paper was written during the meeting "Nonlinear Days in Turin" organized in Turin, Italy, in September 2017. The hosting institution is acknowledged.

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